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VORTICAL INTERACTION OF A SPATIAL LAMINAR BOUNDARY LAYER ON A C--ETC(U)
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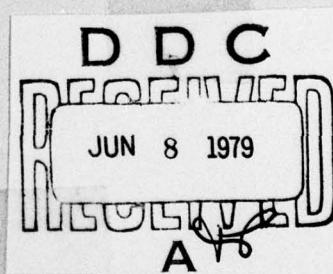
FOREIGN TECHNOLOGY DIVISION



VORTICAL INTERACTION OF A SPATIAL LAMINAR BOUNDARY
LAYER ON A CIRCULAR CONE WITH THE EXTERNAL
(NONVISCOSUS) FLOW AT SUPERSONIC SPEEDS

by

B. M. Bulakh



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А а	А а	A, a	Р р	Р р	R, r
Б б	Б б	B, b	С с	С с	S, s
В в	В в	V, v	Т т	Т т	T, t
Г г	Г г	G, g	У у	У у	U, u
Д д	Д д	D, d	Ф ф	Ф ф	F, f
Е е	Е е	Ye, ye; E, e*	Х х	Х х	Kh, kh
Ж ж	Ж ж	Zh, zh	Ц ц	Ц ц	Ts, ts
З з	З з	Z, z	Ч ч	Ч ч	Ch, ch
И и	И и	I, i	Ш ш	Ш ш	Sh, sh
Й й	Й й	Y, y	Щ щ	Щ щ	Shch, shch
К к	К к	K, k	Ь ъ	Ь ъ	"
Л л	Л л	L, l	Н н	Н н	Y, y
М м	М м	M, m	Ь ъ	Ь ъ	'
Н н	Н н	N, n	Э э	Э э	E, e
О о	О о	O, o	Ю ю	Ю ю	Yu, yu
П п	П п	P, p	Я я	Я я	Ya, ya

*ye initially, after vowels, and after ъ, ъ; e elsewhere.
When written as ё in Russian, transliterate as yё or ё.

RUSSIAN AND ENGLISH TRIGONOMETRIC FUNCTIONS

Russian	English	Russian	English	Russian	English
sin	sin	sh	sinh	arc sh	\sin^{-1}
cos	cos	ch	cosh	arc ch	\cosh^{-1}
tg	tan	th	tanh	arc th	\tanh^{-1}
ctg	cot	cth	coth	arc cth	\coth^{-1}
sec	sec	sch	sech	arc sch	\sech^{-1}
cosec	csc	csch	csch	arc csch	csch^{-1}

Russian	English
rot	curl
lg	log

VORTICAL INTERACTION OF A SPATIAL LAMINAR BOUNDARY LAYER ON A CIRCULAR CONE WITH THE EXTERNAL (NONVISCOSUS) FLOW AT SUPERSONIC SPEEDS

B. M. Bulakh

We have detected the phenomenon of a vortical interaction of a boundary layer and external flow with moderate supersonic speeds. It was shown that for a circular cone the phenomenon of vortical interaction is described by a self-modeling solution.

We are examining the problem of a boundary layer on a circular cone, streamlined by a uniform supersonic flow of gas under the angle of attack α . The author established earlier (in a work published at the MZhG) that in a system of spherical coordinates r , θ, φ , where axis $\theta=0$ coincides with the axis of symmetry of the cone, the behavior of density ρ , component of speed u, w , in the direction of an increase, respectively, of r and φ in the vicinity of the surface of the cone is determined by the formula

$$f = f_0(\tau) + f_1(\tau)(\theta - \theta_*)^B + \dots \quad (1)$$

where $f = \rho, u, w; f_0, f_1$ - some functions of $\varphi; \theta_*$ - angle of the partial opening of the cone; B - the constant which in the majority of

cases is less than a unit; a dotted line designates the members of a higher order of smallness with respect to ϵ than those given. For cases of $B \ll 1$ the nonviscous flow near the surface of the cone whirls powerfully, which renders a significant influence on the boundary layer. If Re is the Reynolds number of the problem, $\epsilon = (Re)^{-\frac{1}{2}}$, then in the system of coordinates where s is read off from the top of the cone along its generatrix, n - along the normal to the surface of the cone, the solution to the problem on the boundary layer has the form:

$$\left. \begin{aligned} u &= u_1(\zeta, \varphi) + \epsilon^B s^{-B} u_2(\zeta, \varphi) + o(\epsilon^B), \\ w &= w_1(\zeta, \varphi) + \epsilon^B s^{-B} w_2(\zeta, \varphi) + o(\epsilon^B), \\ v &= \frac{\epsilon}{\sqrt{s}} [V_1(\zeta, \varphi) + \epsilon^B s^{-B} V_2(\zeta, \varphi) + o(\epsilon^B)], \\ \zeta &= \zeta_1(\zeta, \varphi) + \epsilon^B s^{-B} \zeta_2(\zeta, \varphi) + o(\epsilon^B), \\ T &= t_1(\zeta, \varphi) + \epsilon^B s^{-B} t_2(\zeta, \varphi) + o(\epsilon^B), \\ p &= p_1(\varphi) + o(\epsilon^B), \end{aligned} \right\} \quad (2)$$

here u , v , w - components of the speed of particles of gas in the direction of an increase, respectively, of s , n , φ ; ρ - density and pressure, T - absolute temperature: $\rho := \frac{N}{s}$, $N = \pi \epsilon^{-1}$.

The members with subscript "1" in formulas (2) give the known (self-modeling) solution to the problem of boundary layer; members with subscript "2" occur as a result of the vortical interaction of the external flow and the boundary layer. Equations for u_2, v_2, w_2, t_2, V_2 are analogous to the equations for u_1, v_1, w_1, t_1, V_1 ; boundary conditions are different; therefore they are given thusly:

$$\begin{aligned} \zeta &= 0, \quad u_1 = V_1 = w_1 = 0, \quad t_1 = T_\infty, \\ \zeta &\rightarrow \infty, \quad u_1 \rightarrow u_0(\varphi), \quad w_1 \rightarrow w_0(\varphi), \quad t_1 \rightarrow t_0(\varphi); \\ \zeta &= 0, \quad u_2 = V_2 = w_2 = 0, \quad t_2 = 0, \\ \zeta &\rightarrow \infty, \quad u_2 \sim A_1(\varphi) \epsilon^B, \quad w_2 \sim A_2(\varphi) \epsilon^B, \quad t_2 \sim A_3(\varphi) \epsilon^B. \end{aligned}$$

Here functions $u_1, w_0, r_1, A_1, A_2, A_3$ are determined from the solution to the problem for the external, nonviscous flow; T_w - temperature of the surface of the cone. An indirect evaluation of the value of a check for the results of the normal theory of the boundary layer due to the vortical interaction can be done by means of evaluating the multiplier ε^B in formulas (2).

For example, for the Mach number of an undisturbed flow $M=5$ $\theta_c = 20^\circ$, $\alpha = 10^\circ$, $B \approx 0.33$; for values $M=7$ $\theta_c = 30^\circ$, $\alpha = 5^\circ$, $B \approx 0.075$; if we take $Re = 10^6$, then $\varepsilon = 10^{-3}$ and for the modes given above the flows of the cone are, respectively:

$$\varepsilon^B \approx 0.1 \text{ and } \varepsilon^B \approx 0.6.$$

Consequently, we can anticipate that the corresponding checks are achieved in a number of cases to tens of percents.

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